

“Many black hole” space-times

Piotr T. Chruściel

based on work in collaboration with
Erwann Delay and Rafe Mazzeo

Cargèse, Corsica, August 2002

Brill-Lindquist initial data: the space-metric at time $t = 0$ takes the form

$$g = \psi^4(dx^2 + dy^2 + dz^2) ,$$

with

$$\psi = 1 + \sum_{i=1}^I \frac{m_i}{2|\vec{x} - \vec{x}_i|} .$$

A special case: the space-part of the Schwarzschild metric centred at \vec{x}_0 with mass m :

$$g = \left(1 + \frac{m}{2|\vec{x} - \vec{x}_0|} \right)^4 \delta ,$$

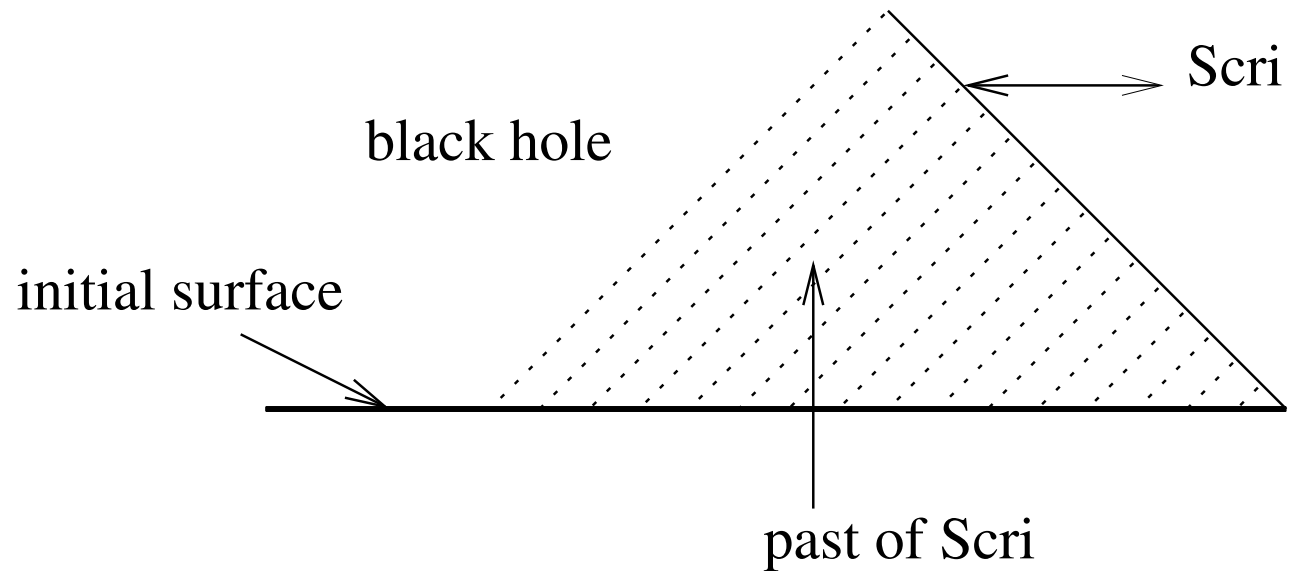
where δ is the Euclidean metric.

Question: How differentiable is the associated Scri?

What is a black hole in a space-time (\mathcal{M}, g) ?

\mathcal{B} — black hole region

$$\mathcal{B} := \mathcal{M} \setminus J^-(\mathcal{I}^+)$$



\mathcal{I}^+ — future null infinity Scri,

cf. Friedrich's and Galloway's lectures

Brill-Lindquist initial data: the space-metric at time $t = 0$ takes the form

$$g = \psi^4(dx^2 + dy^2 + dz^2) ,$$

with

$$\psi = 1 + \sum_{i=1}^I \frac{m_i}{2|\vec{x} - \vec{x}_i|} .$$

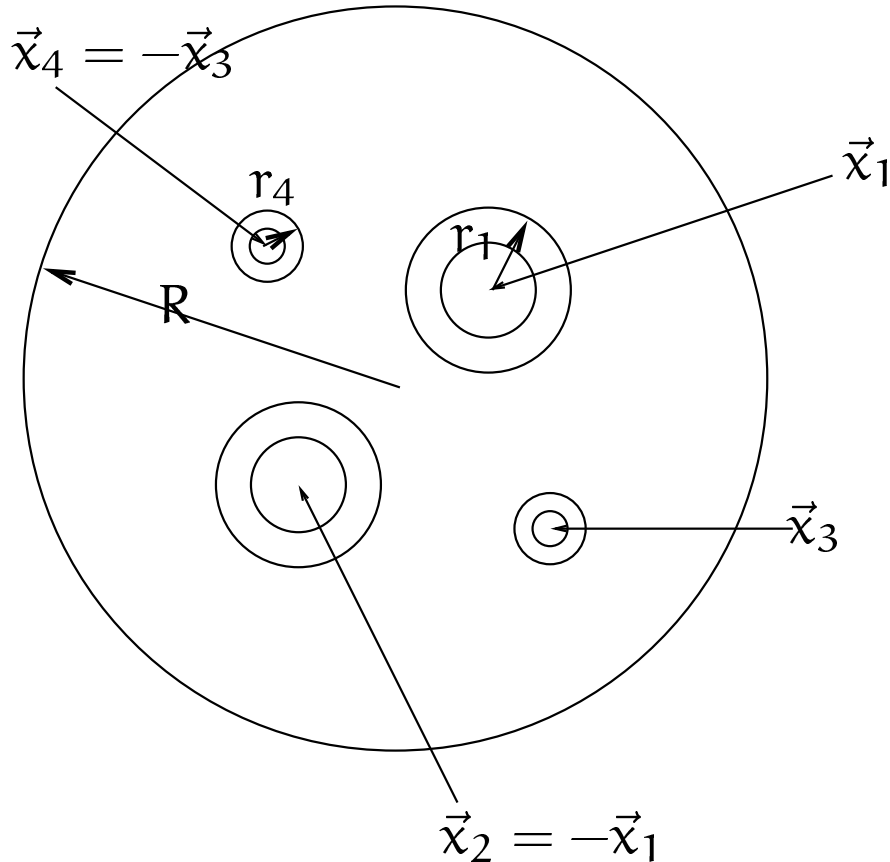
A special case: the space-part of the Schwarzschild metric centred at \vec{x}_0 with mass m :

$$g = \left(1 + \frac{m}{2|\vec{x} - \vec{x}_0|} \right)^4 \delta ,$$

where δ is the Euclidean metric.

Question: How differentiable is the associated Scri?

The “many Schwarzschild metrics” (PTC, E.Delay, CQG, gr-qc/0203053) a solution exists when $0 \leq m_i \leq \delta$, δ small enough



Free parameters: R and
 (\vec{x}_1, r_1, m_1) , (\vec{x}_3, r_3, m_3)

$$m_2 = m_1, \quad r_2 = r_1$$

$$m_4 = m_3, \quad r_4 = r_3$$

inner small circles are Schwarzschild horizons; metric *exactly*
Schwarzschild within the outer small circles and outside the big circle

Recall:

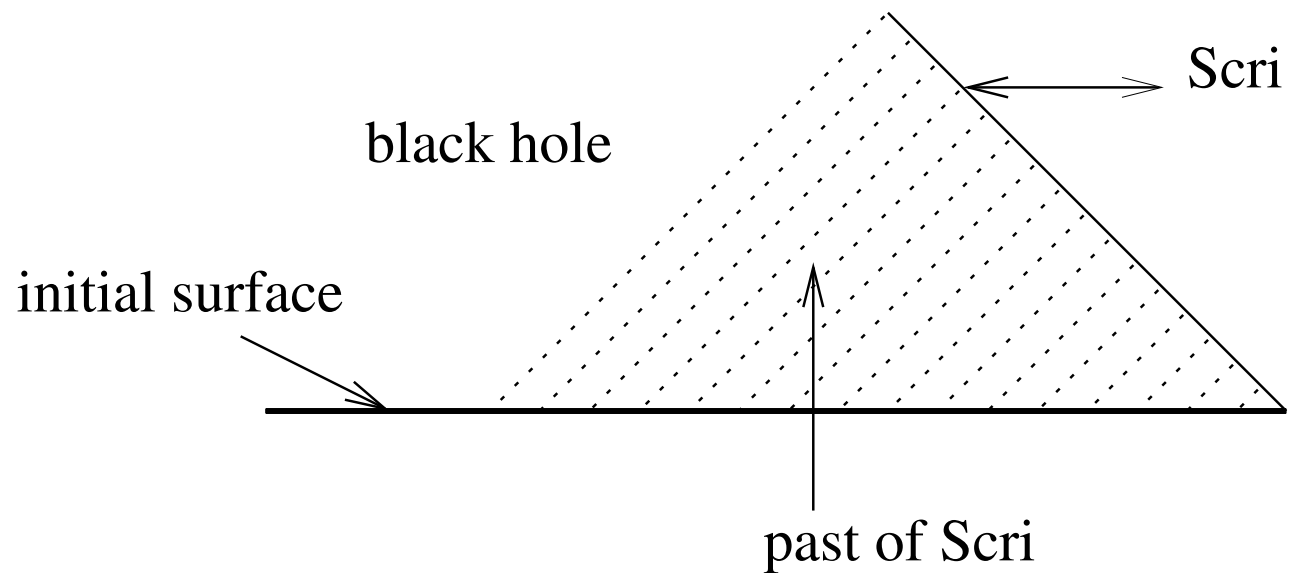
- for time-symmetric initial data ($K_{ij} = 0$):
apparent horizon = outermost minimal surface
- **If** there exists a sufficiently regular conformal completion Scri , then
apparent horizon \implies black hole region

Theorem: (PTC, R.Mazzeo, in preparation) If the m_i 's $i = 1 \dots I$, are small enough, then the *apparent horizon* (= outermost minimal surface) has at least (precisely?) I connected components.

What is a black hole in a space-time (\mathcal{M}, g) ?

\mathcal{B} — black hole region

$$\mathcal{B} := \mathcal{M} \setminus J^-(\mathcal{I}^+)$$



\mathcal{I}^+ — future null infinity Scri,

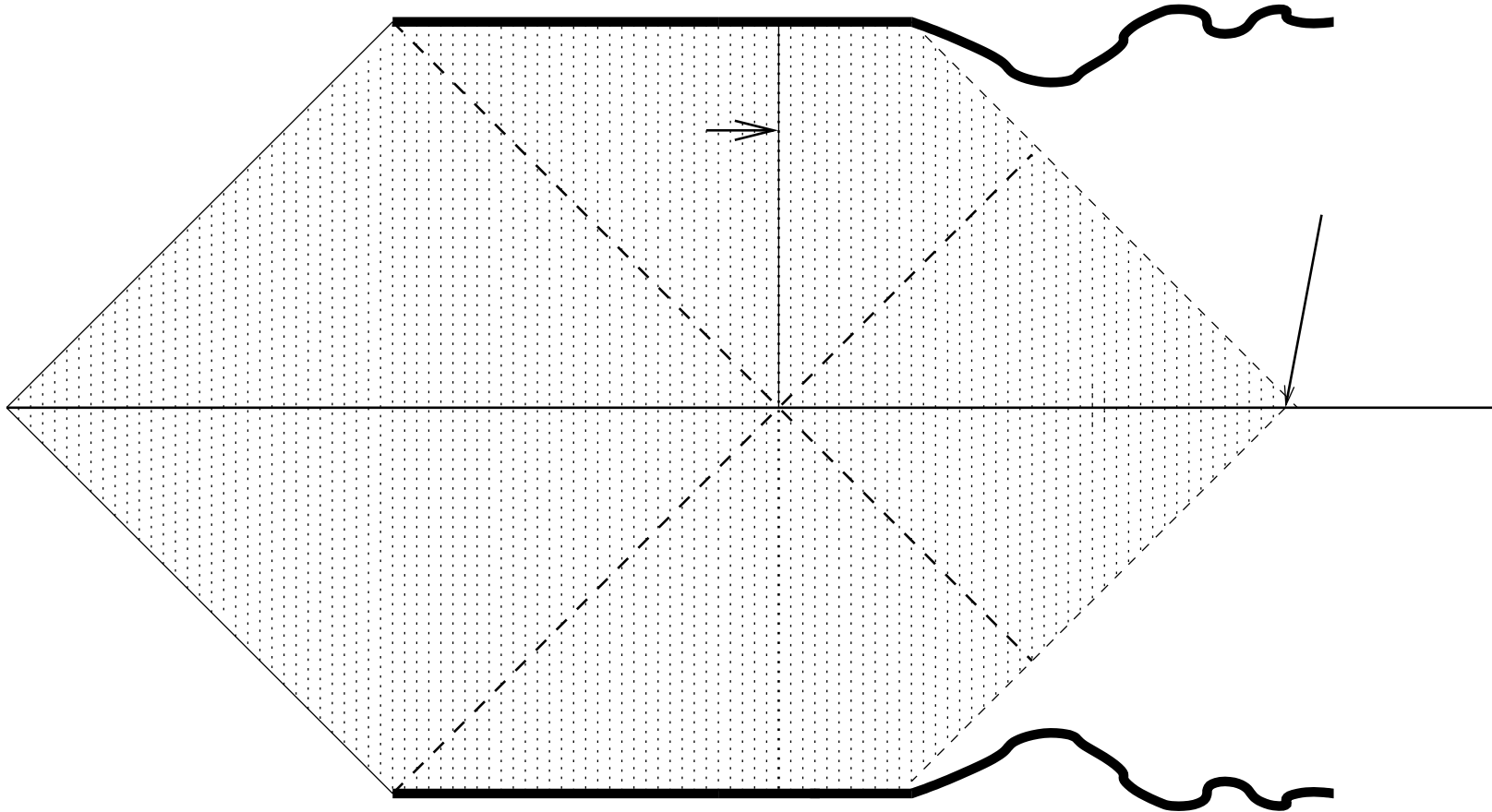
cf. Friedrich's and Galloway's lectures

Recall:

- for time-symmetric initial data ($K_{ij} = 0$):
apparent horizon = outermost minimal surface
- **If** there exists a sufficiently regular conformal completion Scri , then
apparent horizon \implies black hole region

Theorem: (PTC, R.Mazzeo, in preparation) If the m_i 's $i = 1 \dots I$, are small enough, then the *apparent horizon* (= outermost minimal surface) has at least (precisely?) I connected components.

What about event horizons? The resulting space-times are “small”:

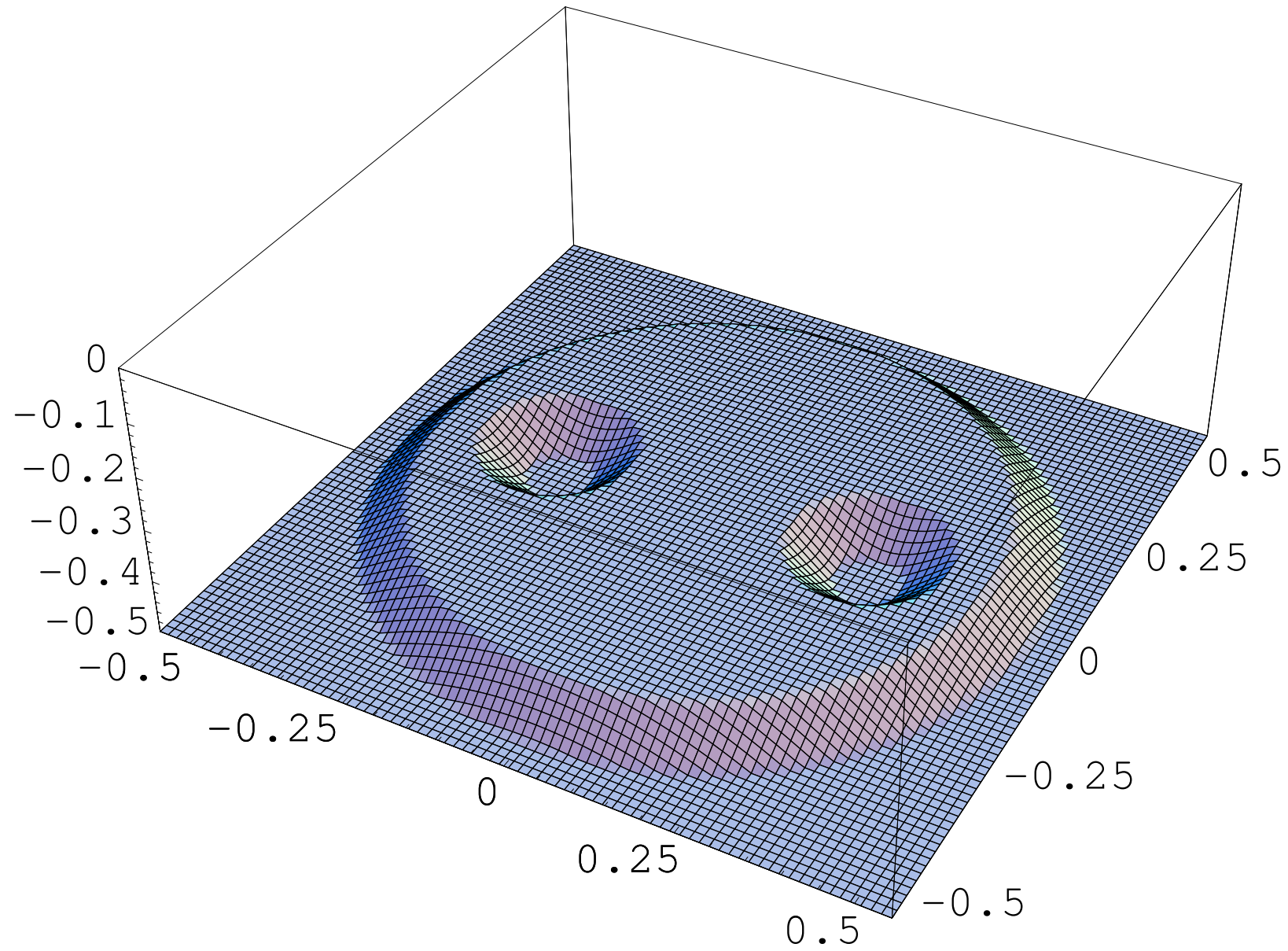


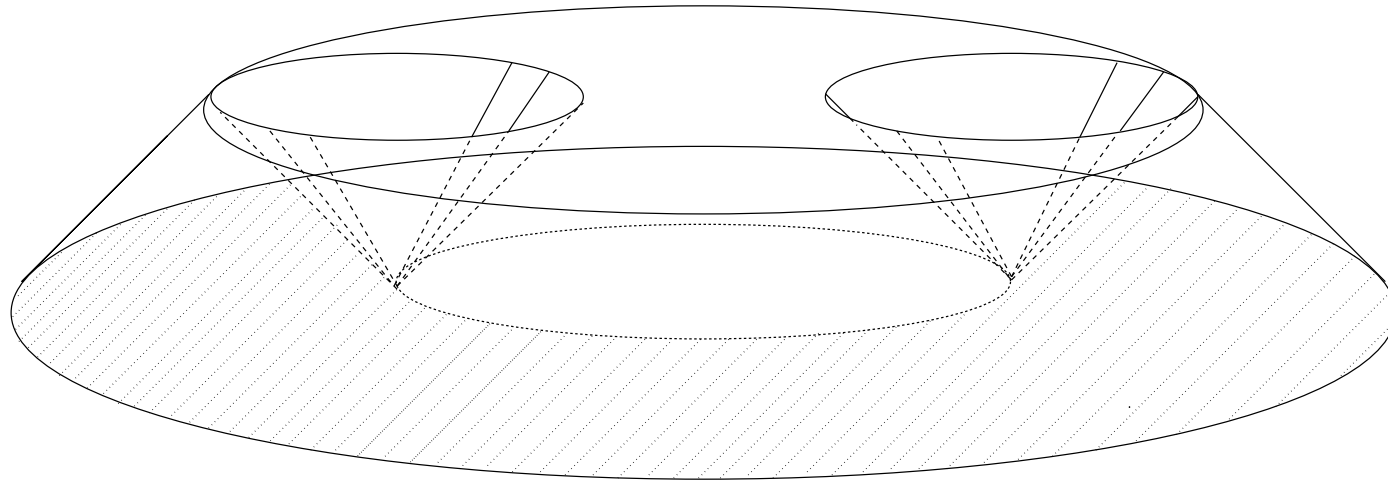
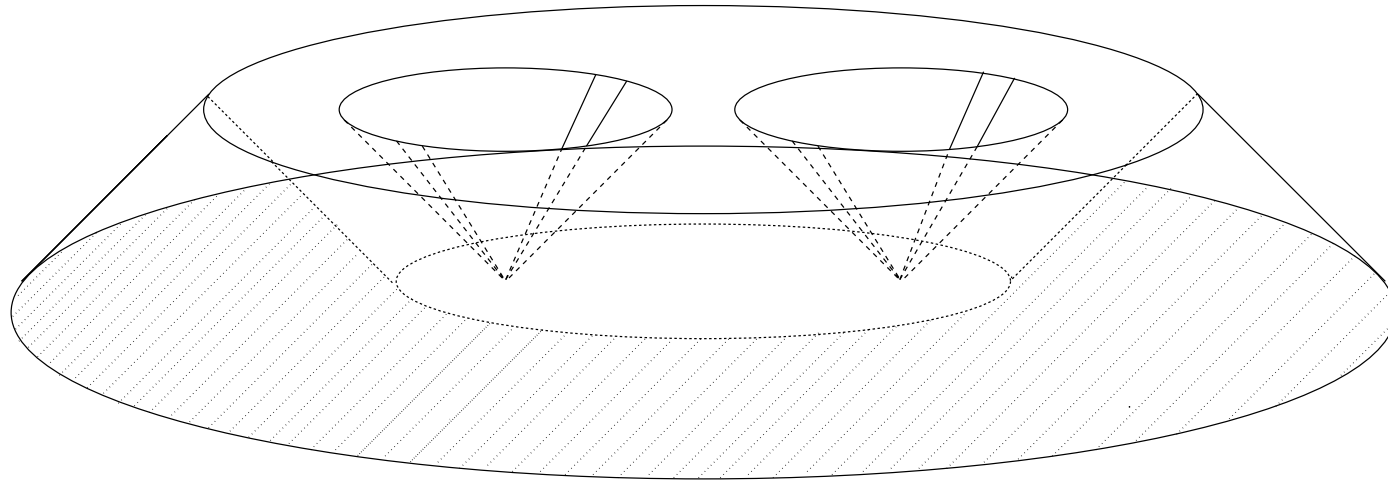
Theorem (PTC, R.Mazzeo, in preparation)

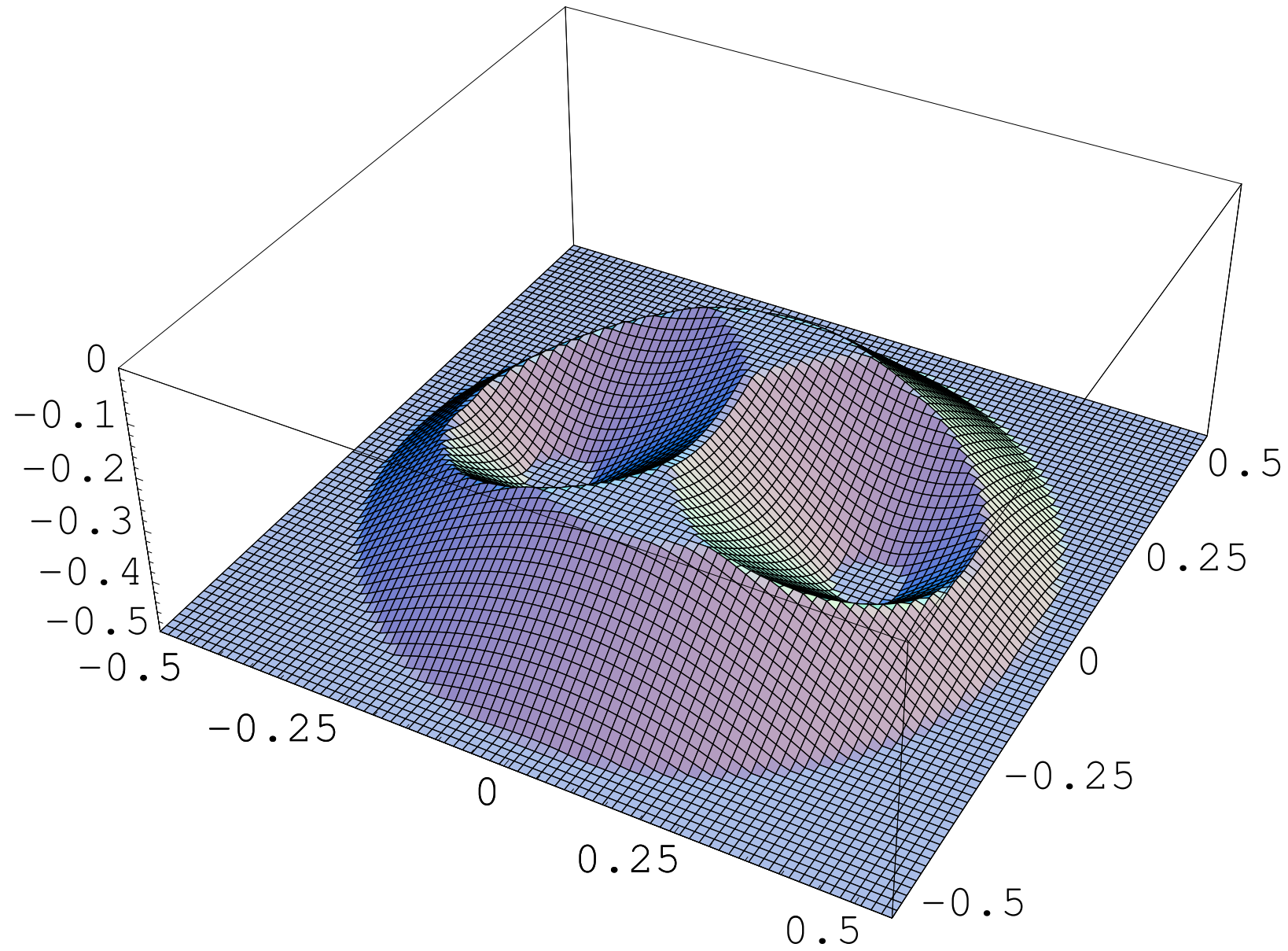
There exist configurations with *at least* I connected components of the intersection of the event horizon with the initial data hypersurface:

- two black holes, m_i small enough
- all black holes aligned, m_i small enough.

Why is there a problem? A possible global structure:







A simple way of obtaining (a piece of) the Minkowskian Scri:

$$\{x^0 > 0, \eta_{\alpha\beta} x^\alpha x^\beta < 0\} \ni x^\mu \rightarrow y^\mu = \frac{x^\mu}{\eta_{\alpha\beta} x^\alpha x^\beta} \in \{y^0 < 0, \eta_{\alpha\beta} y^\alpha y^\beta < 0\}.$$

Here $\eta_{\mu\nu}$ is the Minkowski metric.

$y^\mu = 0$ is a point representing the "timelike future infinity", usually denoted by i^+

